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The Bromine Nuclear Quadrupole Resonance of Two Modifications of Indium Tribromide

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The bromine nuclear quadrupole resonance study of indium tribromide has revealed that there are at least two modifications in its solid state. The high-temperature modification shows the two 81 Br resonance lines of 85.985 and 86.205 MHz, with the relative intensity of 1: 2 at 24.5°C. This modification was stable unless the temperature was lowered to $\sim 15^{\circ}$ C. The low-temperature modification shows one 81 Br resonance line, 86.915 MHz, at -196° C and was stable up to $\sim 30^{\circ}$ C. The directions of the z axes of the field gradient were determined by studying the Zeeman effect on each resonance line of the two modifications at room temperature. For the two resonance lines of the high-temperature modification, one and two directions, which had the asymmetry parameters of 0.362 and 0.400 respectively, were observed. The resonance line of the low-temperature modification showed three directions with the asymmetry parameter of 0.387. The crystal structures of the high- and low-temperature modifications were, therefore, concluded to be monoclinic and hexagonal respectively, from the numbers of resonance lines and of the nonequivalent directions of the z axes. The quadrupole coupling constants and asymmetry parameters were analyzed by means of the Townes-Dailey method.

Indium trihalides, InX₃ (X: Cl, Br, I), exist as dimers in the gaseous state,¹⁾ but their crystal structures are not always the same. According to the X-ray diffraction analysis,²⁾ the crystal of InCl₃ has a layer structure

in which every indium atom is octahedrally surrounded by chlorine atoms, and its space group is C2/m as in AlCl₃.³⁾ On the other hand, the crystal of InI₃ is made up of dimer molecules joined by iodine bridges,⁴⁾

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as in the gaseous state. As for InBr₃, the X-ray analysis has not yet been reported, as far as we know.

The results of the nuclear quadrupole resonance (NQR) of bromine in InBr₃ were reported previously by Ludwig⁵ and Barnes et al.,⁶ but they are not completely consistent. Barnes et al. studied extensively the NQR on the trihalides of Group IIIb elements and found that all the NQR spectra consist of a singlet on the lower-frequency side and a doublet on the higher-frequency side.^{6,7} The ⁸¹Br NQR spectrum of InBr₃ consists of the singlet of 86.912 MHz and the doublet of 106.974 and 107.428 MHz at -196°C. The former is ascribed to the bridging bromine atoms, while the latter is assinged to the terminal ones.

On the other hand, Greenwood et al. have reported the infrared and Raman spectra of InX₃.8) According to their results, the spectra of InBr₃ are more similar to those of InCl₃ than those of InI₃; that is, the crystal of InBr₃ has a layer structure. This is contrary to the NQR results presented above.

We re-examined the ⁸¹Br NQR of InBr₃ and, furthermore, observed the Zeeman effect on the NQR. We obtained different spectra from those reported previously; our results can be satisfactorily explained by the layer structure suggested by Greenwood *et al.*

Also, the study of the temperature dependences of the NQR frequencies revealed the existence of at least two modifications in InBr₃ crystals. The high-temperature modification shows two resonance lines, while the low-temperature modification shows only one resonance line. The measurements of the NQR Zeeman effect were made on these two modifications.

Experimental

The sample of InBr₃ was purchased from Katayama Chemical Industries Co., Ltd., and was purified by sublimation. The single crystal necessary for the study of the Zeeman effect was grown by the Bridgman-Stockbarger method.

The NQR spectrometer used in these experiments was a super-regenerative oscillator with frequency modulation; the absorption lines were observed on an oscilloscope. The resonance frequencies were determined with a VHF signal generator and a frequency counter, TR-5578, from the Takeda Riken Industry Co., Ltd. The Zeeman effect was examined by means of the zero-splitting cone method used in a low-magnetic field. The magnetic field of about 200 G was applied by means of a Helmholtz coil, and the patterns of the zero-splitting loci were determined at room temperature. The temperature dependences of the resonance frequencies were measured by using a thermostat from the Takara Kogyo Co. After the sample had been cooled or heated at a fixed temperature for 30—60 min, the resonance frequencies were determined.

Results and Discussion

The ⁸¹Br NQR frequencies of InBr₃ observed in our experiment are presented in Table 1. Those for ⁷⁹Br were observed at frequencies corresponding to the frequency ratio (1.197) of ν (⁷⁹Br)/ ν (⁸¹Br). The crystals of InBr₃ commercially obtained or purified by sublimation always show two closely-spaced resonance lines at room temperature. The intensity ratio of the higher line (ν_{β_2}) to the lower (ν_{β_1}) is approximately 2:1. These results indicate that there are two kinds of nonequivalent bromine atoms in the crystal and that their abundance ratio is 2:1. On cooling, this modification is apparently transformed to the other one, which shows only one resonance line. Tentatively we will refer to the former as the β -modification and to the latter as the α -modification.

Table 1. ⁸¹Br nuclear quadrupole resonance frequencies in the modifications of indium tribromide

Modification		Resonace frequency ^{a)} (MHz)			
		−196°C	24.5°C	49°C	
α	v_{α}	86.915	85.962		
β	v_{β_1}		85.985(1) ^{b)}	85.867(1)	
	v_{β_2}		86.205(2)	86.074(2)	

a) Experimental error is ± 0.01 MHz.

The resonance frequencies in Table 1 are close to that of the bridging bromine atom reported by Barnes et al.⁶⁾ Especially, the resonance frequency of the α -modification (ν_{α}) at -196° C is consistent within the limits of experimental error. We scanned carefully the frequency regions around 105 and 125 MHz, where the resonance frequencies of the terminal ⁸¹Br and ⁷⁹Br atoms in the dimer have been reported by Barnes et al.,⁶⁾ but we could not detect any resonance line.⁹⁾ Therefore, we concluded that neither of the modifications observed here is composed of dimers.

The temperature dependences of the resonance frequencies were measured between -48 and 49° C, as is shown in Fig. 1. The transition point is hard to determine because of the large hysteresis. In the course of cooling the single crystal of the β -modification, an absorption line of the α -modification appeared at $\sim 15^{\circ}$ C. On warming, the signal of the α -modification was observed up to $\sim 30^{\circ}$ C. Using the powder sample, the region of the hysteresis spread as may be seen in Fig. 1. Furthermore, in the temperature range from -6 to -16° C, the absorption lines of both the α - and β -modifications were observed.

The β -Modification of Indium Tribromide. The single crystal of InBr₃ initially obtained had the β -modification, because it was crystallized from the melt. The patterns of zero-splitting were obtained on each resonance line. For the bromine atom (I=3/2), the

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⁷⁾ R. G. Barnes and S. L. Segel, *ibid.*, **25**, 180 (1956). S. L. Segel and R. G. Barnes, *ibid.*, **25**, 578 (1956).

⁸⁾ N. N. Greenwood, D. J. Prince, and B. P. Straughan, J. Chem. Soc., A, 1968, 1964.

b) The value in parenthesis is the relative intensity.

⁹⁾ The sensitivity of the spectrometer was confirmed by the detection of ^{127}I resonance line of InI_3 for 125 MHz region. The ^{79}Br resonance lines of $AlBr_3$ and also, of $InBr_3$ give the enough sensitivity for 105 MHz region.

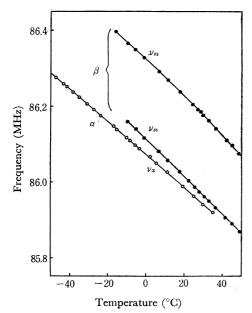


Fig. 1. Temperature dependences of ⁸¹Br nuclear quadrupole resonance frequencies in indium tribromide.

zero-splitting occurs under the following conditions:¹⁰⁾
$$\sin^2 \theta = 2/(3 - \eta \cos 2\varphi),$$
 (1)

where the orientation of the magnetic field is expressed by the polar coordinates (θ,φ) with respect to the coordinates of the principal axes (x, y, and z axes) of the field gradients at the resonant nuclei and where η is an asymmetry parameter of the field gradient. In this way, a cone about the z axis can be drawn by θ and φ which satisfies Eq. (1). The axes corresponding to $\varphi=0$ and $\varphi=\pi/2$ are shown by the x and y axes respectively. Only one zero-splitting pattern was obtained for ν_{β_1} , whereas two patterns were observed for ν_{β_2} . The experimental results are shown in Fig. 2, where Z_1 , Z_{2A} , and Z_{2B} are the directions of the z axes at the bromine atoms contributing to ν_{β_1} and ν_{β_2} . From the number of the z axes of ν_{β_1} and ν_{β_2} , the symmetry of the β -modification must be lower than the tetragonal,

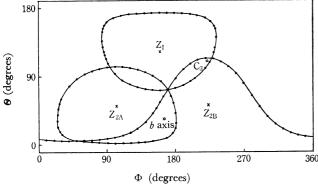


Fig. 2. Zero-splitting patterns of the Zeeman effect of the β -modification of indium tribromide. Θ and Φ are polar and azimuthal angles in the coordinates fixed to the sample respectively.

but not the triclinic.¹¹⁾ However, as a symmetry higher than the monoclinic is ruled out by the requirement for the angles between the z axes, it can be concluded that the β -modification belongs to the monoclinic space group. The asymmetry parameters of the field gradient were determined from Eq. (1) by the use of the least-squares method. As may be seen in Table 2, large values of η were obtained.

Table 2. Asymmetry parameters and quadrupole coupling constants for ⁸¹Br in indium tribromide at room temperature

Modification		η	$e^2Qq/h({ m MHz})$	Up	
α	(v_{α})	0.387 ± 0.005	167.79	0.2609	
β	$(\nu_{\boldsymbol{\beta_1}})$	0.362 ± 0.002	168.33	0.2618	
	$(\nu_{\boldsymbol{\beta_2}})$	0.400 ± 0.007	167.99	0.2612	

Greenwood et al. suggested that the crystal structure of $InBr_3$ is similar to that of $InCl_3$ or $AlCl_3$.⁸⁾ In the case of the $AlCl_3$ -type structure, there are two non-equivalent halogen atoms which form slightly different bridges. Therefore, the low-resonance frequencies and the large asymmetry parameters may be attributed to the bridging bromine atoms in the β -modification. Also, it is reasonable, in view of this structure, that the β -modification shows two resonance lines with the intensity ratio of 1:2 and that the zero-splitting patterns observed are one for ν_{β_1} and two for ν_{β_2} . Table 3

Table 3. Angles between the z axes and the b axis in the β -modification of indium tribromide

	$Z_1 - Z_2$	$Z_{2\mathrm{A}} - Z_{2\mathrm{B}}$	$Z_1 - b$ axis	Z_2-b axis
NQR	92°08′	88°16′	90°	44°08′
X-Ray ^{a)}	90°27′	88°26′	90°	44°13′

a) These values were calculated from the atomic positions of AlCl₃. (Ref. 3).

shows the angles between the z axes and the b axis of the monoclinic lattice as determined by the Zeeman effect. The values in the second row are calculated from the atomic positions of AlCl₃,3) because the data of X-ray analysis are not available for InBr₃. In this calculation, the z axis was assumed to be directed perpendicularly to the plane including M-X-M, as in the case of a bridging halogen atom of the dimer molecule. 12,13) The results of the NQR and those of the X-ray analysis are in good agreement, though the latter are derived for a related compound, AlCl₃. Accordingly, it seems most likely that the β -modification is isomorphous with AlCl₃. This conclusion is further supported by the fact that Zeeman patterns similar to those in Fig. 2 have been found in ErCl₃, which has an AlCl₃-type structure, by Carlson and Adams.14)

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¹²⁾ G. E. Peterson and P. M. Bridenbaugh, J. Chem. Phys., 51, 238 (1969).

¹³⁾ T. Okuda, H. Terao, O. Ege, and H. Negita, *ibid.*, **52**, 5489 (1970).

¹⁴⁾ E. H. Carlson and H. S. Adams, ibid., 51, 388 (1969).

The \(\alpha\)-Modification of Indium Tribromide. When the single crystal of the β -modification mounted on the apparatus was cooled to about -50° C by means of a cold stream of nitrogen gas from liquid nitrogen, it was transformed to the single crystal of the α -modification. The measurements of the zero-splitting on this modification show three different directions of the z axes, indicating that the crystal of the \alpha-modification has a threefold axis. 15) In Fig. 2, C3 indicates the orientation of this threefold axis of the α-modification at the same coordinates of the β -modification. The compound that possesses three nonequivalent directions of the z axes belongs, in general, to the hexagonal (point group C₃ or C_{3i}) system or, in special cases, to another hexagonal or cubic system. In this case, however, the cubic system, which has three z axes orthogonal with one another,11) can be excluded because the observed angle between the z axes is 92°13′.

The zero-splitting patterns of the α -modification are very similar to those of the β -modification. The directions of the z axes of the α -modification deviate from the relevant axes of the β -modification to a small degree. It can, therefore, be expected that the crystal structure of the α-modification is slightly different from that of the β -modification; that is, the crystal of the α-modification may be composed of a polymeric layer lattice like AlCl₃. In view of these results, it is likely to be of the BiI₃-type, the structure of which is closely related to the AlCl₃-type¹⁶⁾ and which has the space group R3.17) The 127I NQR of BiI3 indicates only a resonance line and a large asymmetry parameter of the field gradient.¹⁸⁾ The transition may occur between monoclinic and hexagonal systems, as in the case of CrCl₃ reported by Morosin and Narath.¹⁹⁾

The crystal of $InBr_3$ is easily cloven along the plane perpendicular to the threefold axis, which may be parallel to the c axis of the hexagonal unit cell. The b axis of the monoclinic lattice of the β -modification was perpendicular to the threefold axis of the α -modification; that is, it was parallel to the cleavage plane.

It is interesting to notice that InBr₃ has a higher melting point (436°C) than those of the other Group

IIIb trihalides which form dimers in the crystal (e.g., 210°C for InI₃²⁰) and that the ratio of the ionic radii for In³+/Br⁻ is 0.42. This value is slightly larger than the lower limiting ratio, 0.414, for the octahedral coordination of the metal atom. In view of these points, InBr₃ can be expected to form an infinite layer lattice, as is proposed by this experiment.

Bond Character. The analyses of the Zeeman effect of InBr₃ reveal a close relation between the α -and β -modifications. Such a result as above can also be seen from Tables 1 and 2, *i.e.*, the resonance frequencies and the asymmetry parameters in both modifications are almost the same.

The observed x axis of the field gradient at the bromine atom is directed nearly along the bisector of In-Br-In and the y axis is parallel to the In-In direction. The z axis is perpendicular to the plane including In-Br-In, as has been explained before. From these coordinates of the principal field gradients, we can construct the hybrid orbitals of the bromine atom to bind with two indium atoms. Using the Townes and Dailey theory and appropriate approximations, the following two equations are found to hold for the NQR parameters: 21)

$$Up(1+\eta/3) = 2 - a, (3)$$

where a is the electron population in the bonding orbital directed to the indium atom and where α is the interbond angle, \angle In-Br-In. Up is the number of unbalanced p electrons and can be determined from the observed quadrupole coupling constant, $e^2Qq_{\rm obsd}$:

$$hv = 1/2e^2Qq_{\text{obsd}}(1+\eta^2/3)^{1/2},$$
 (4)

$$Up = e^2 Q q_{\text{obsd}} / e^2 Q q_{\text{atom}} \tag{5}$$

where $e^2Qq_{\rm atom}/h$ is 643.03 MHz.²²⁾

Substituting the observed values of η and Up into Eqs. (2) and (3), we obtain $a{=}1.7$ as the mean value, which gives a negative charge of 0.4e on the bromine atom. This ionic character is smaller than the 55% expected from the electronegativity difference between indium and bromine. This discrepancy may be due to the crystal field exerted by the neighboring nuclei. The angles of In–Br–In are deduced to be $96^{\circ}56' \sim 97^{\circ}40'$. These values are somewhat larger than the relevant angles, $94^{\circ}20'$ and $95^{\circ}30'$ for AlCl₃.

¹⁵⁾ The NQR Zeeman analysis always introduces the center of symmetry. It is impossible to distinguish between the threefold rotation axis and the threefold inversion axis on the basis of the Zeeman analysis. (Ref. 11).

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